Variable Eddington Factor Acceleration of Thermal Radiative Transfer on Curved Meshes

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- Thermal Radiative Transfer: describes conservation and transfer of energy between photons and matter
- 6+1 dimensional phase space \Rightarrow dominates memory and runtime
- LDRD investigating high order FEM TRT on curved meshes for coupling to BLAST hydrodynamics code
- Goal: develop an acceleration scheme to improve iterative convergence









• Steady state, one-group, Linear Boltzmann Equation with isotropic scattering and source:

$$\hat{\Omega} \cdot \nabla \psi + \sigma_t \psi = \frac{\sigma_s}{4\pi} \int \psi \, \mathrm{d}\Omega + \frac{Q}{4\pi}$$
Streaming Collision Scattering Source

- Discrete Ordinates (S_N) angular discretization

$$\hat{\Omega} \cdot \nabla \psi_d + \sigma_t \psi_d = \frac{\sigma_s}{4\pi} \sum w_{d'} \psi_{d'} + \frac{Q}{4\pi} \,, \quad d = 1, 2, \dots, N_{\text{angles}}$$

where $\psi_d(x) = \psi(x, \hat{\Omega}_d)$

• $N_{\rm angles}$ coupled equations \Rightarrow prohibitively expensive to solve

• Decouple by lagging the scattering term

$$\hat{\Omega} \cdot \nabla \psi_d^{\ell+1} + \sigma_t \psi_d^{\ell+1} = \frac{\sigma_s}{4\pi} \sum w_{d'} \psi_{d'}^{\ell} + \frac{Q}{4\pi}$$
Known from previous iteration

- $\rightarrow N_{\rm angles}$ independent equations but need to solve iteratively
- Slow convergence in highly scattering systems



Need a preconditioner/accelerator for practical applications

Variable Eddington Factor Equations

• Take first two angular moments of transport equation:

$$\nabla \cdot \vec{J} + \sigma_a \phi = Q$$
$$\nabla \cdot (\boldsymbol{E}\phi) + \sigma_t \vec{J} = 0$$

where

$$\phi = \int \psi \, \mathrm{d}\Omega \,, \quad \vec{J} = \int \hat{\Omega} \, \psi \, \mathrm{d}\Omega$$

and

$$oldsymbol{E} = rac{\int \hat{\Omega} \otimes \hat{\Omega} \, \psi \, \mathrm{d}\Omega}{\int \psi \, \mathrm{d}\Omega}$$

- In 3D: 4 equations for 13 unknowns
- More angular moments \rightarrow more unknowns
- Need ψ (the solution) to have closure
- Historically: invent a closure similar to flux limited diffusion

Close with transport information from previous iteration

Solve

$$\hat{\Omega} \cdot \nabla \psi^{\ell+1/2} + \sigma_t \psi^{\ell+1/2} = \frac{\sigma_s}{4\pi} \phi^\ell + \frac{Q}{4\pi}$$

for $\psi^{\ell+1/2}$

• Compute Eddington tensor:

$$\boldsymbol{E}_{ij}^{\ell+1/2} = \frac{\sum w_d \hat{\Omega}_i^{(d)} \hat{\Omega}_j^{(d)} \psi_d^{\ell+1/2}}{\sum w_d \psi_d^{\ell+1/2}}$$

- Solve VEF equations for updated scalar flux $\phi^{\ell+1}$

$$\nabla \cdot \vec{J}^{\ell+1} + \sigma_a \phi^{\ell+1} = Q ,$$
$$\nabla \cdot \left(\boldsymbol{E}^{\ell+1/2} \phi^{\ell+1} \right) + \sigma_t \vec{J}^{\ell+1} = 0$$

- Update scattering term with VEF solution
- Stop if $\|\phi^{\ell+1} \phi^{\ell}\| < tol$



- Eddington tensor converges faster than the scalar flux
 - Angular flux weighted average of $\hat\Omega\otimes\hat\Omega\Rightarrow$ depends on angular shape not magnitude
 - ψ converges quickly in angular shape
- VEF compensates lagging of scattering term in Source Iteration

Mixed Finite Element Discretization

VEF discretization does not need to match transport! Mixed Finite Element:

- \vec{J} with $H^{1,d}$ finite elements (vector lagrange)
- ϕ with L^2 finite elements (discontinuous)

Multiply zeroth moment by ϕ basis function, u, and integrate:

$$\int u\nabla \cdot \vec{J_h} \,\mathrm{d}V + \int \sigma_a u \phi_h \,\mathrm{d}V = \int uQ \,\mathrm{d}V$$

Multiply first moment by \vec{J} basis function, \vec{v} , and integrate tensor term by parts:

$$\int \phi_h \boldsymbol{E} : \nabla \vec{v} \, \mathrm{d}V - \int \sigma_t \vec{v} \cdot \vec{J}_h \, \mathrm{d}V = \int_{\partial V} \boldsymbol{E} \hat{n} \cdot \vec{v} \phi_h^{\mathsf{BC}} \, \mathrm{d}S$$

 $\nabla \cdot \vec{J} + \sigma_a \phi = Q$

 $\nabla \cdot (\boldsymbol{E}\phi) + \sigma_t \vec{J} = 0$

Matrix form:

$$\begin{bmatrix} \mathbf{M}_{a} & \mathbf{G} \\ \mathbf{H} & -\mathbf{M}_{t} \end{bmatrix} \begin{bmatrix} \underline{\phi} \\ \underline{J} \end{bmatrix} = \begin{bmatrix} \underline{Q} \\ \underline{B} \end{bmatrix}$$

where

$$\begin{split} \mathbf{M}_{a,ij} &= \int \sigma_a u_i u_j \, \mathrm{d}V \,, \qquad \mathbf{G}_{ij} = \int u_i \nabla \cdot \vec{v}_j \, \mathrm{d}V \,, \\ \mathbf{H}_{ij} &= \int u_j \boldsymbol{E} : \nabla \vec{v}_i \, \mathrm{d}V \,, \qquad \mathbf{M}_{t,ij} = \int \sigma_t \vec{v}_i \cdot \vec{v}_j \, \mathrm{d}V \\ \underline{Q}_i &= \int u_i Q \, \mathrm{d}V \,, \qquad \underline{B}_i = \int_{\partial V} \boldsymbol{E} \hat{n} \cdot \vec{v}_i \phi_h^{\mathsf{BC}} \, \mathrm{d}S \end{split}$$

 $\mathbf{H}\neq\mathbf{G}^{T}$ due to presence of Eddington Tensor \Rightarrow non-symmetric Saddle Point Problem

Solve with Schur Complement

 $\begin{bmatrix} \mathbf{M}_a & \mathbf{G} \\ \mathbf{H} & -\mathbf{M}_t \end{bmatrix} \begin{bmatrix} \underline{\phi} \\ \underline{J} \end{bmatrix} = \begin{bmatrix} \underline{Q} \\ \underline{B} \end{bmatrix}$

• \mathbf{M}_a is block diagonal \Rightarrow easily inverted!

$$\mathbf{M}_{a}\underline{\phi} + \mathbf{G}\underline{J} = \underline{Q}$$
$$\Rightarrow \underline{\phi} = \mathbf{M}_{a}^{-1} \left[\underline{Q} - \mathbf{G}\underline{J} \right]$$

$$\begin{split} \mathbf{H}\underline{\phi} - \mathbf{M}_t \underline{J} &= \underline{B} \\ \Rightarrow -\underbrace{\left[\mathbf{H}\mathbf{M}_a^{-1}\mathbf{G} + \mathbf{M}_t\right]}_{\mathsf{Schur Complement} = \mathbf{S}} \underline{J} = \underline{B} - \mathbf{H}\mathbf{M}_a^{-1}\underline{Q} \end{split}$$

- Assemble and solve at every iteration \Rightarrow want iterative solver
- Schur Complement is non-symmetric and has been difficult to solve iteratively (need to find a preconditioner)

Geometry: 2D box $[0,1]\times [0,1]$

 $\sigma_t = 5 \, \mathrm{cm}^{-1} \Rightarrow 5 \mathrm{~mfp}$ thick

Set the fixed source to force the solution to "chopped sines"

$$\phi = \sin\left(\pi\frac{x+\alpha}{1+2\alpha}\right)\sin\left(\pi\frac{y+\alpha}{1+2\alpha}\right)$$

 $\alpha \neq 0$ allows testing inflow boundary conditions

Provides known solution to compare numerical solution to

 S_4 angular quadrature





Fig. 1: Convergence for unaccelerated Source Iteration.

VEF Accelerates Source Iteration



Fig. 2: Convergence for VEF acceleration.

VEF Accelerates Source Iteration





Fig. 2: Convergence for VEF acceleration.

VEF algorithm maintains order of accuracy



VEF accelerates as well as traditional methods



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- Developed a mixed finite element discretization for Eddington equations + a robust Saddle Point solver
- Implemented VEF acceleration in LDRD code
- Verified order of accuracy and acceleration properties

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