## High Order Discontinuous Galerkin Discrete Ordinates Solver

NE 255 Final Project Presentation

Samuel S. Olivier November 29, 2018



## Motivation

- Finite Element Method
  - Strong mathematical foundation (weak forms, approximation theory)  $\Rightarrow$  good error analysis
  - Unstructured grids  $\Rightarrow$  match complex geometries easily
- High Order FEM
  - More accurate per unknown (in smooth problems)
  - Use fewer more accurate but more expensive elements  $\Rightarrow$  less communication
  - $\bullet~\mbox{Curved}$  meshes  $\Rightarrow$  even better for complex geometry
- Standard continuous FEM not stable for hyperbolic PDEs  $\Rightarrow$  Discontinuous Galerkin FEM
- Goal: implement DGFEM for  $S_N$  in 1D slab geometry





• Steady state, one-group, Linear Boltzmann Equation with isotropic scattering and source:



•  $N_{\text{angles}}$  coupled equations  $\Rightarrow$  prohibitively large system of equations

#### Source Iteration decouples in angle

• Decouple by lagging the scattering term

$$\mu_d \frac{\mathrm{d}\psi_d^{\ell+1}}{\mathrm{d}x} + \Sigma_t \psi_d^{\ell+1} = \frac{\Sigma_s}{2} \sum w_{d'} \psi_{d'}^{\ell} + \frac{Q}{2}$$
Known from previous iteration

 $\rightarrow N_{\rm angles}$  independent equations but need to solve iteratively

• At every iteration, need to solve

$$\mu_d \frac{\mathrm{d}\psi_d}{\mathrm{d}x} + \Sigma_t \psi_d = S$$

for every angle



## **Discontinuous Galerkin Derivation**

- Derive for "master" element on  $\xi \in [-1,1]$
- Approximate with linear combination of interpolating, nodal basis functions in each element

$$\psi_d(\xi) = \sum_j \psi_{d,j} b_j(\xi), \quad b_j \in P_k(\xi)$$

- DGFEM procedure:
  - Plug in FEM approximation for  $\psi_d$
  - Multiply equation by  $b_i(\xi)$
  - Integrate over the element
  - Use integration by parts to get a boundary term
  - Apply upwinding to boundary term to uniquely define the element edges
  - Map to a "real" element



## Discontinuous Galerkin Derivation (cont.)

• Multiply by a "test" function  $b_i(\xi)$  and integrate over the "master" element  $\xi \in [-1, 1]$ :

$$\sum_{j} \psi_{d,j} \int_{-1}^{1} \mu_{d} b_{i} \frac{\mathrm{d}b_{j}}{\mathrm{d}x} + \Sigma_{t} b_{i} b_{j} \,\mathrm{d}\xi = \int_{-1}^{1} b_{i} S \,\mathrm{d}\xi$$

• Integrate by parts:

$$[\mu_d b_i \psi_d]_{-1}^{+1} + \sum_j \psi_{d,j} \int_{-1}^1 -\mu_d \frac{\mathrm{d}b_i}{\mathrm{d}x} b_j + \Sigma_t b_i b_j \,\mathrm{d}\xi = \int_{-1}^1 b_i S \,\mathrm{d}\xi$$

Note

$$[\mu_d b_i \psi_d]_{-1}^{+1} = \mu_d b_i(1) \psi_d(1) - \mu_d b_i(-1) \psi_d(-1)$$

• How to evaluate  $\psi_d(\pm 1)$ ?



#### Use upwinding to uniquely define the edges

• Upwinding = search backward along particle path

$$\mu_d > 0: \quad \psi_d(-1) = \psi_{\text{in}}, \quad \psi_d(+1) = \sum_j b_j(+1)\psi_{d,j}$$
$$\mu_d < 0: \quad \psi_d(-1) = \sum_j b_j(-1)\psi_{d,j}, \quad \psi_d(+1) = \psi_d$$

$$\mu_d < 0: \quad \psi_d(-1) = \sum_j b_j(-1)\psi_{d,j}, \quad \psi_d(+1) = \psi_{\text{in}}$$

where  $\psi_{\rm in}$  is incoming flux known from the boundary condition or an upwind element



#### Map to physical space

• Convert master to physical space with isoparametric transformation

$$x(\xi) = \sum_{i} b_i(\xi) x_i, \quad J(\xi) = \frac{\mathrm{d}x}{\mathrm{d}\xi} = \sum_{i} \frac{\mathrm{d}b_i}{\mathrm{d}\xi} x_i$$

where  $x_i$  are the node locations in physical space

• Can integrate any function of x in "reference" space  $\xi \in [-1, 1]$  with

$$\int_{-1}^{1} f(x(\xi)) J(\xi) \,\mathrm{d}\xi$$

• Evaluate numerically with Gauss Quadrature



## **Matrix Form**

Can rewrite in matrix form:

$$\left[\mu_d \mathbf{G} + \mathbf{M}\right] \vec{\psi}_d = \vec{S} + \vec{B} \,,$$

where

$$\mathbf{G}_{ij} = \begin{cases} \psi_{d,j}b_i \Big|_{\xi=+1} - \int_{-1}^{1} \frac{\mathrm{d}b_i}{\mathrm{d}\xi} b_j \,\mathrm{d}\xi \,, & \mu_d > 0\\ -\psi_{d,j}b_i \Big|_{\xi=-1} - \int_{-1}^{1} \frac{\mathrm{d}b_i}{\mathrm{d}\xi} b_j \,\mathrm{d}\xi \,, & \mu_d < 0 \end{cases}, \quad \mathbf{M}_{ij} = \int_{-1}^{1} \Sigma_t b_i b_j J(\xi) \,\mathrm{d}\xi \,,$$

$$\vec{B}_{i} = \begin{cases} \mu_{d}\psi_{\mathrm{in}}v_{i}\Big|_{\xi=-1}, & \mu_{d} > 0\\ -\mu_{d}\psi_{\mathrm{in}}v_{i}\Big|_{\xi=+1}, & \mu_{d} < 0 \end{cases}, \qquad \qquad \vec{S}_{i} = \int_{-1}^{1} b_{i}SJ(\xi) \,\mathrm{d}\xi$$

Matrices are  $(p+1) \times (p+1)$ 

### **DGFEM Sweeps**

- Start with  $\psi_{\rm in}={\rm known}$  boundary condition
- Solve for the  $\vec{\psi}_d$  in the first element with:

$$\vec{\psi}_d = \left[\mu_d \mathbf{G} + \mathbf{M}\right]^{-1} \left[\vec{S} + \vec{B}\right]$$

by computing the inverse (the system is  $p + 1 \times p + 1 \Rightarrow$  small)

• Compute the incoming angular flux for the next element as the outgoing from the previous element

$$\psi_{\rm in} = \sum_j b_j(1)\psi_{d,j}$$

- Repeatedly solve using the outgoing flux from the previous element until the right boundary is reached
- Same process for negative angles but sweep from right to left

- Implemented in C++ ( $\sim$  3k LOC!)
- Calls LAPACK for dense linear algebra (matrix inversions, matrix vector products, etc)
- Uses Lua scripts for parameter input i.e. number of angles, number of elements, cross sections, source
- CMake build system
- Outputs in Vislt's curve file format for visualization

• Choose known solution, plug into transport equation, and solve for the source

$$\psi_d(x) = \frac{1}{2}\sin(\pi x) \Rightarrow \phi(x) = \sin(\pi x)$$
$$\Rightarrow Q_d = \frac{1}{2}[\mu_d \pi \cos(\pi x) + \Sigma_a \sin(\pi x)]$$

- Provides analytic solution to compare to
- Expect

$$E = Ch^{p+1}$$

in the  $L^2$  norm

• All calculations were  $\mathsf{S}_{32}$ 

#### MMS test verifies arbitrary order of accuracy



#### High order is more accurate per unknown



Number of Unknowns

## Solution for Inhomogeneous Source

$$\Sigma_t = 1 \,\mathrm{cm}^{-1}\,, \quad \Sigma_s = 0.8 \,\mathrm{cm}^{-1}\,, \quad Q = \begin{cases} 1\,, & x < 4 \\ 0\,, & x \ge 4 \end{cases}$$



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- Implemented DGFEM  $S_N$  solver in C++
- Verified code with order of accuracy MMS tests
- Showed that high order is more accurate per unknown

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# Questions?